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# How relevant is charge-particle transport in global models?

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## Popular choice to study plasma chemistry

- adopt a **spatial averaged description**,  
hence involving little computational effort
  - allow describing in detail the plasma chemistry  
in complex gas mixtures
  - should include transport effects,  
especially at low to intermediate pressures
- usually involve the coupled solution of  
**Chemistry solver** (to solve the “plasma chemistry”)  
**Boltzmann solver** (to describe the “electron kinetics”)


# Transport in global models

The (average) rate balance equation for particle  $s$

$$\frac{\partial n_s}{\partial t} \simeq S_s - \nu_{\text{transp}_s} n_s$$

Average loss-frequency  
due to transport

?


$$\frac{1}{V} \int_{\vec{r}} \vec{\nabla} \cdot \vec{\Gamma}_s d^3r$$

**Focus on charged-particle transport:  
how relevant in global models ?**

# Formulating charged-particle transport

## Defining charged-particle fluxes: challenges

- various working conditions
  - single-/multi- ion scenarios
  - electropositive (EP) / electronegative (EN) plasmas
  - low-pressure (LP) / high-pressure (HP)
- spatial-averaged (algebraic) description
  - simplifications
  - resort to spatially-resolved simulations
  - to define abacuses for  $\nu_{\text{transp}_s}$

- **Transport models in global models**

  - Ambipolar-based models

  - h-factor models

- **Simulations and results**

  - DC oxygen plasma

    - at low-to-intermediate pressures

  - MW helium plasma

    - at low-to-intermediate + atmospheric pressure

- **Final remarks**

Alves and Tejero-del-Caz, PSST (accepted, 2023)

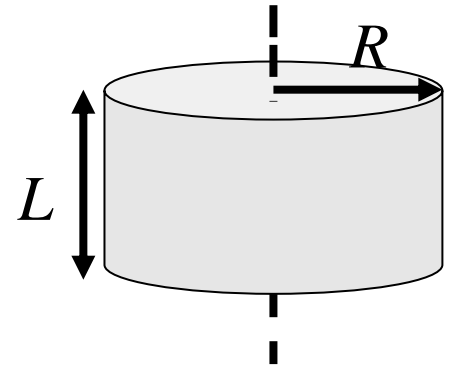
A stylized graphic of a globe, composed of several concentric, curved lines in shades of gray and white, creating a sense of depth and movement. The lines are thicker on the left and become thinner and more numerous towards the right, suggesting a 3D effect.

# **Transport models in global models**

# Ambipolar-based transport models

The (ambipolar) flux for ion species  $i$

$$\vec{\Gamma}_i = -D_{\text{amb}i} \vec{\nabla} n_i$$

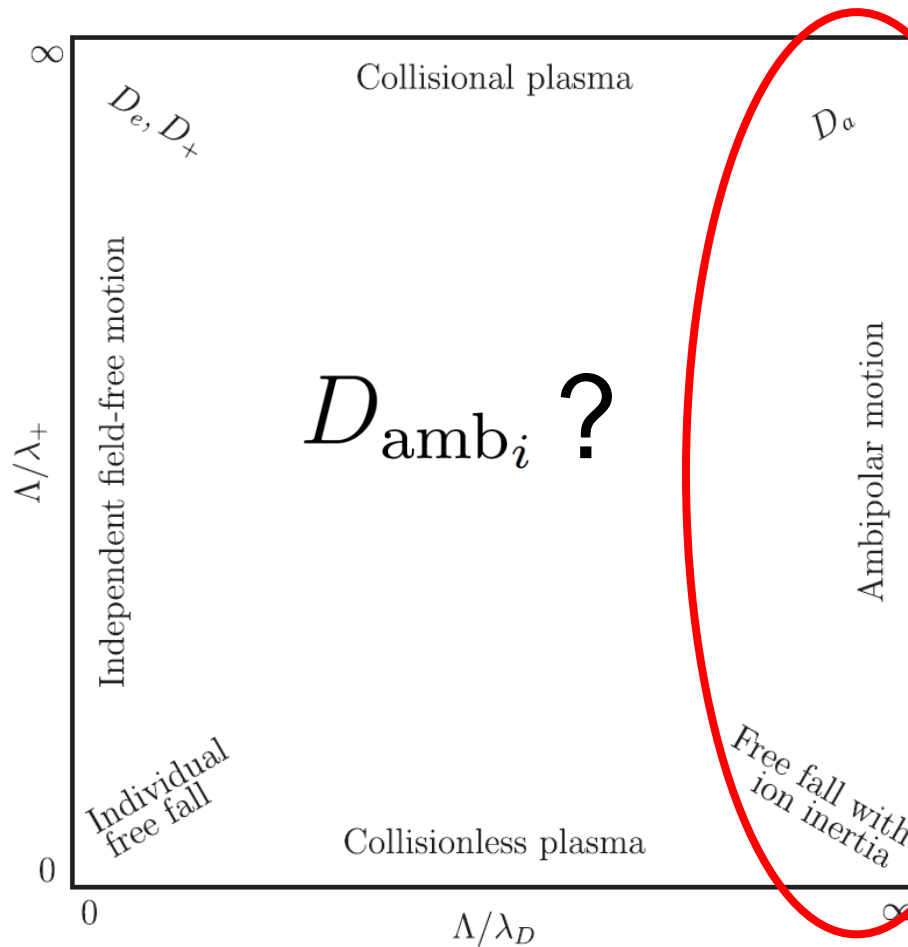


➔

$$\left\{ \begin{array}{l} \nu_{\text{transp}i}^{\text{amb}} = \frac{D_{\text{amb}i}}{\Lambda^2} \\ \Lambda^2 = \left[ \left( \frac{\pi}{L} \right)^2 + \left( \frac{2.405}{R} \right)^2 \right]^{-1} \end{array} \right.$$

# Effective ambipolar diffusion model

unified theory for dense plasmas



**High-pressure:  
classical ambipolar diffusion**

**Dense plasmas**

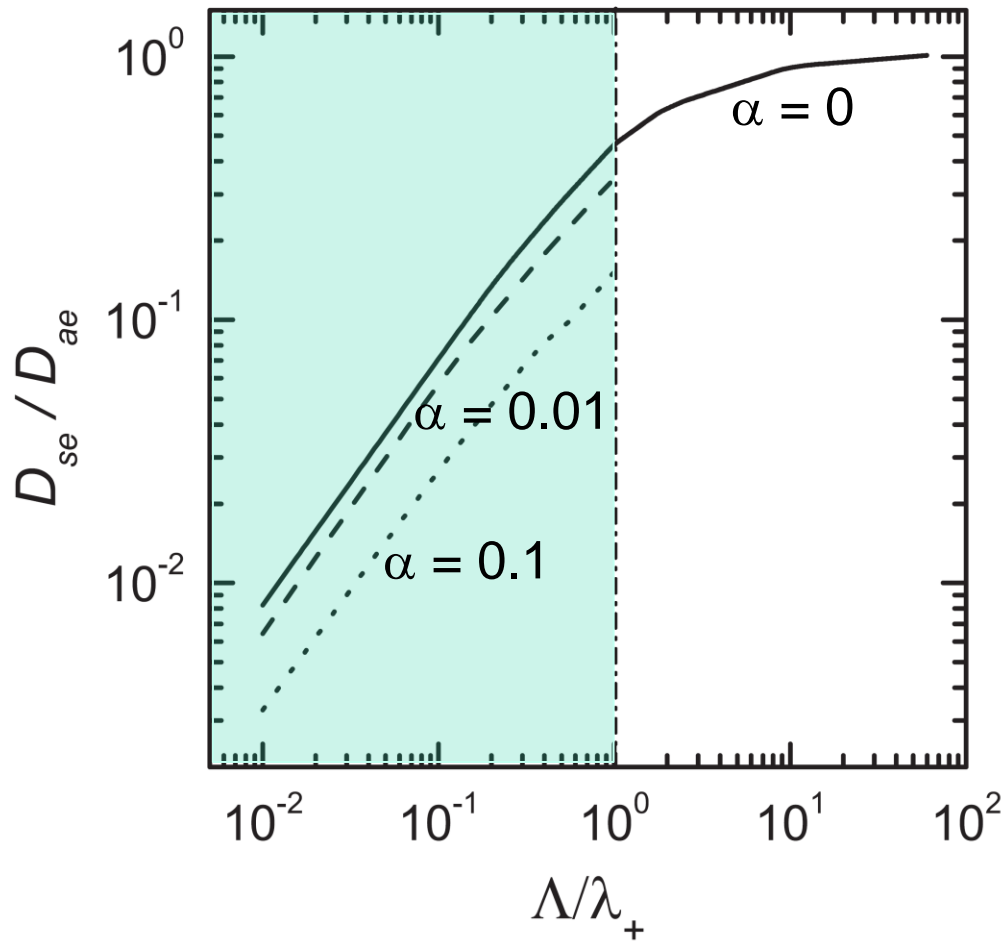
$$\frac{\Lambda}{\lambda_D} \gg 1$$

**Low-pressure:  
free-fall limit**

Coche *et al.*, JPD (2016)  
Phelps, *J. Res. NIST* (1990)



# The effective ambipolar diffusion coefficient



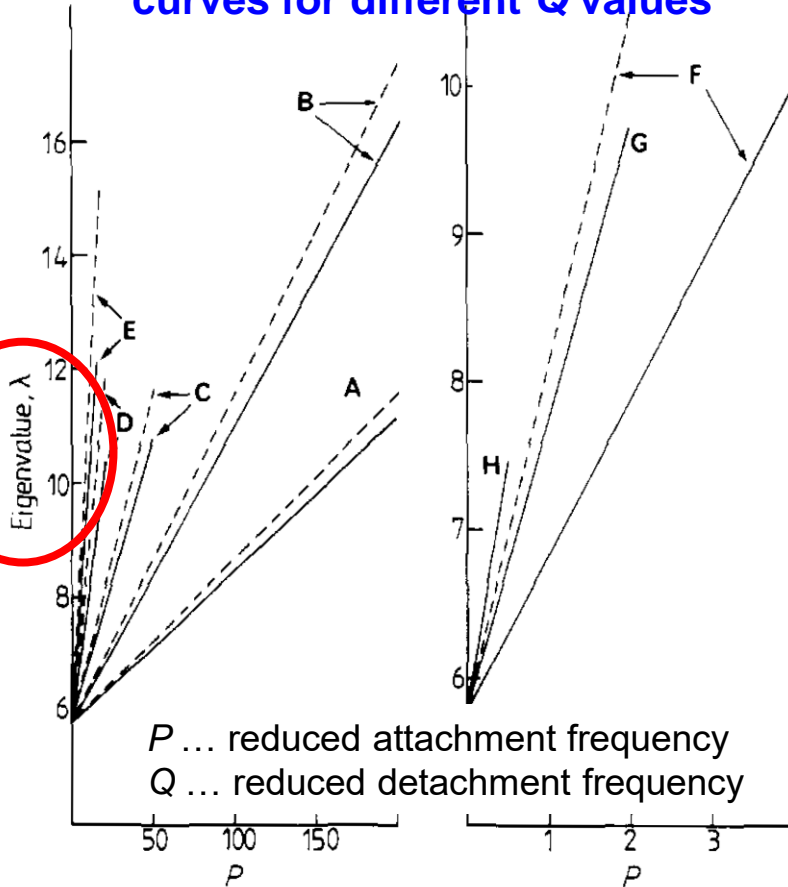
Coche *et al.*, JPD (2016)

## Theory valid for

- Multiple positive ions with similar transport parameters with similar density profiles  
$$\vec{\nabla} n_i / n_i \simeq \vec{\nabla} n_j / n_j$$
- In EP plasmas for all pressures
- In EN plasmas for low pressure for single negative ion with low density ( $\alpha < 0.1$ )

# The eigenvalue diffusion length model

curves for different Q values



P ... reduced attachment frequency  
Q ... reduced detachment frequency

Ferreira *et al.*, JPD (1988)  
Guerra and Loureiro, PSST (1999)

## Theory valid for

- Description of radial transport
- Multiple positive ions with similar creation/destruction rates
- Single negative ion
- Ion creation/destruction due to ionization / recombination attachment / detachment charge transfer

$$\nu_{\text{transp}_i}^{\text{eigen}} = \frac{D_{a_i}}{\Lambda_{\text{eigen}}^2}$$

$$\Lambda_{\text{eigen}} \equiv R \sqrt{\frac{1 + \alpha}{\lambda_{\text{eigen}}}}$$

# The Quantemol Global Model (QGM)

## Ion losses at the wall due to ambipolar transport and thermal losses

$$\nu_{\text{transp}_i}^{\text{QGM}} \approx \frac{1}{\nu_{\text{amb}_i}^{-1} + \nu_{\text{th}_i}^{-1}}$$

$$= \frac{A \Gamma_{\text{wall}_i}}{V n_i} = \frac{A}{V} \frac{D_{a+} \gamma_{r_i}}{\Lambda \gamma_{r_i} + \frac{D_{a+}}{v_{\text{th}_i}}}$$

Tennyson *et al.*, PSST (2022)

$$\frac{1}{V} \int_{\vec{r}} \vec{\nabla} \cdot \vec{\Gamma}_i d^3r = \frac{A}{V} \Gamma_{\text{wall}_i}$$


# The Quantemol Global Model (QGM)

the ambipolar diffusion coefficient

$$D_{a_+} \simeq D_+ \frac{1 + (T_e/T_+)(1 + 2\alpha)}{1 + \alpha(T_e/T_+)}$$

$$D_+ \equiv \frac{\sum_i D_i n_i}{\sum_i n_i}$$

$$D_i = \frac{\pi}{8} \lambda_i v_{thi}$$

$$\frac{1}{\lambda_i} = \sum_j n_j \sigma_{ij}$$


Hard-sphere model / LJ parameters  
Rutherford scattering

## Theory valid for

- Multiple positive and negative ions with similar density profiles

$$\vec{\nabla} n_i / n_i \simeq \vec{\nabla} n_j / n_j$$

- Electron and negative ions in Boltzmann equilibrium with the space-charge potential

- Moderate electronegativities and low ion mobilities

$$\alpha < 1 ; \mu_{+,-} / \mu_e \ll 1$$

Tennyson *et al.*, PSST (2022)

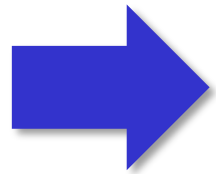
Stoffels *et al.*, *Contribs. Plasma Phys.* (1995)

Rogoff JPD (1985)

# h-factors transport models

## The flux for ion species $i$ at the sheath edge

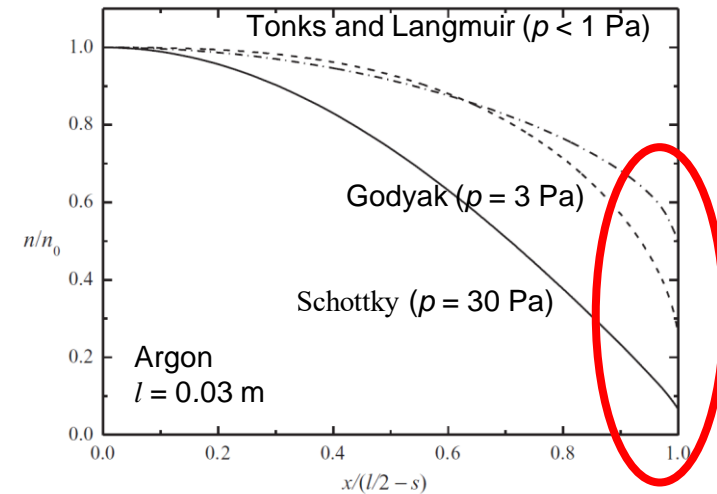
$$\Gamma_{i_{\text{sh}}} = n_{i_{\text{sh}}} u_B$$



$$\nu_{\text{transp}_i}^{\text{h-fact}} \simeq \frac{\Gamma_{i_{\text{sh,axial}}} A_L + \Gamma_{i_{\text{sh,radial}}} A_R}{n_{e0} V}$$
$$= 2u_B \left( \frac{h_{L_i}}{L} + \frac{h_{R_i}}{R} \right)$$

## The h-factors

$$h_{L_i} \equiv \frac{n_{i_{\text{sh,axial}}}}{n_{e0}} ; h_{R_i} \equiv \frac{n_{i_{\text{sh,radial}}}}{n_{e0}}$$

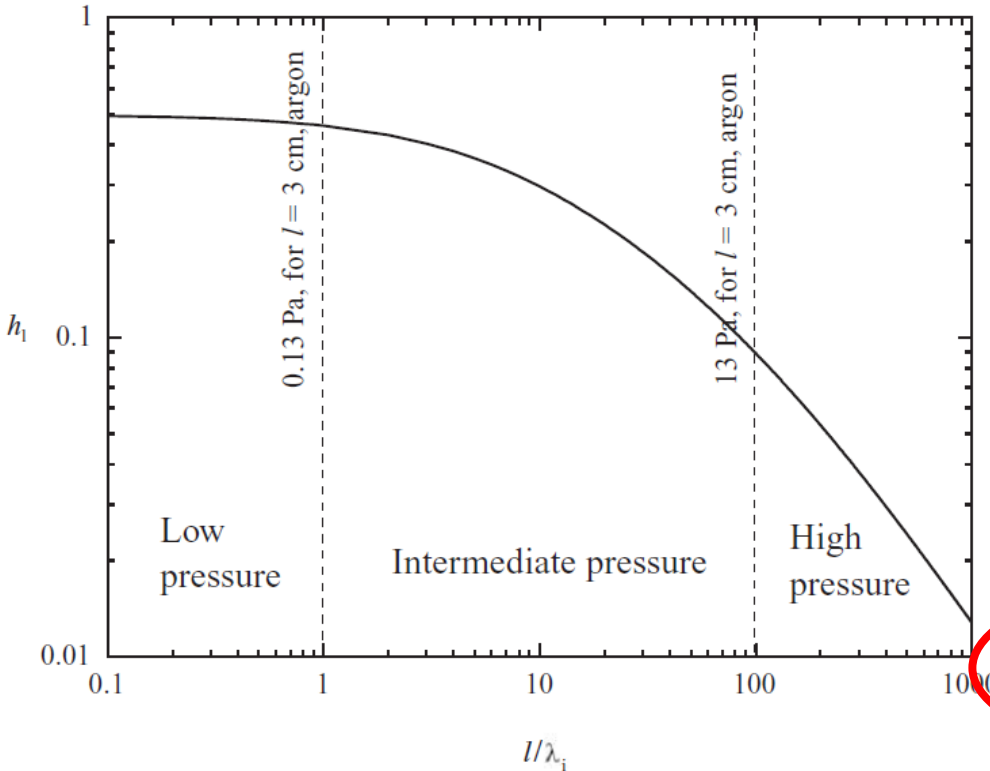


Chabert and Braithwaite (2011)

# h-factors transport models

## the h-factors

Heuristic expressions *stitching together* the various electronegativity / pressure regions



Chabert and Braithwaite (2011)

## Theory valid for

- Multiple positive and negative ions
- Various electronegativity and pressure regimes
  - LP regime, with parabolic-profile EN core and two EP edges;
  - LP regime at high  $\alpha$ , with one-region parabolic EN plasma;
  - HP regime at high  $\alpha$ , with one-region flat-topped EN plasma.

- LP expressions dependent on

$$\lambda_i^{-1} \simeq \sum_j N \sigma_{ij}$$

Thorsteinsson and Gudmundsson, PSST (2010)  
Chabert, PSST (2016)  
Godyak (1986)  
Lichtenberg *et al.*, JAP (1994)  
Lee and Lieberman, JVST (1995)  
Kim *et al.*, JVST (2006)

An abstract graphic consisting of several overlapping circles and curved lines in shades of gray. The circles are of varying sizes and are positioned to create a sense of depth and movement. The lines are curved and appear to be part of the circles or are separate elements that intersect them. The overall effect is a dynamic, layered composition.

# **Simulations and results**

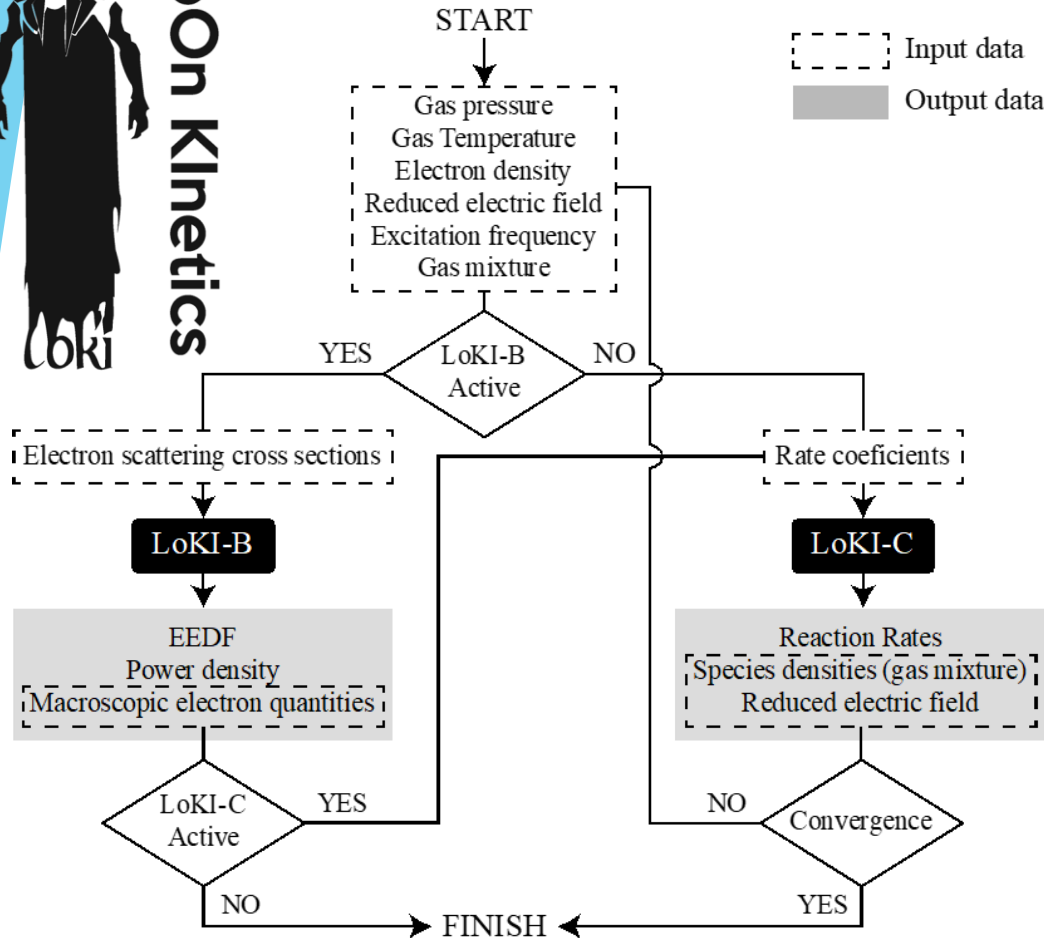
# The LisbOn Knetics (LoKI) simulation tool

(developed under MATLAB®)

**OPEN SOURCE**



LisbOn Knetics



## LoKI-B

<https://github.com/IST-Lisbon/LoKI>

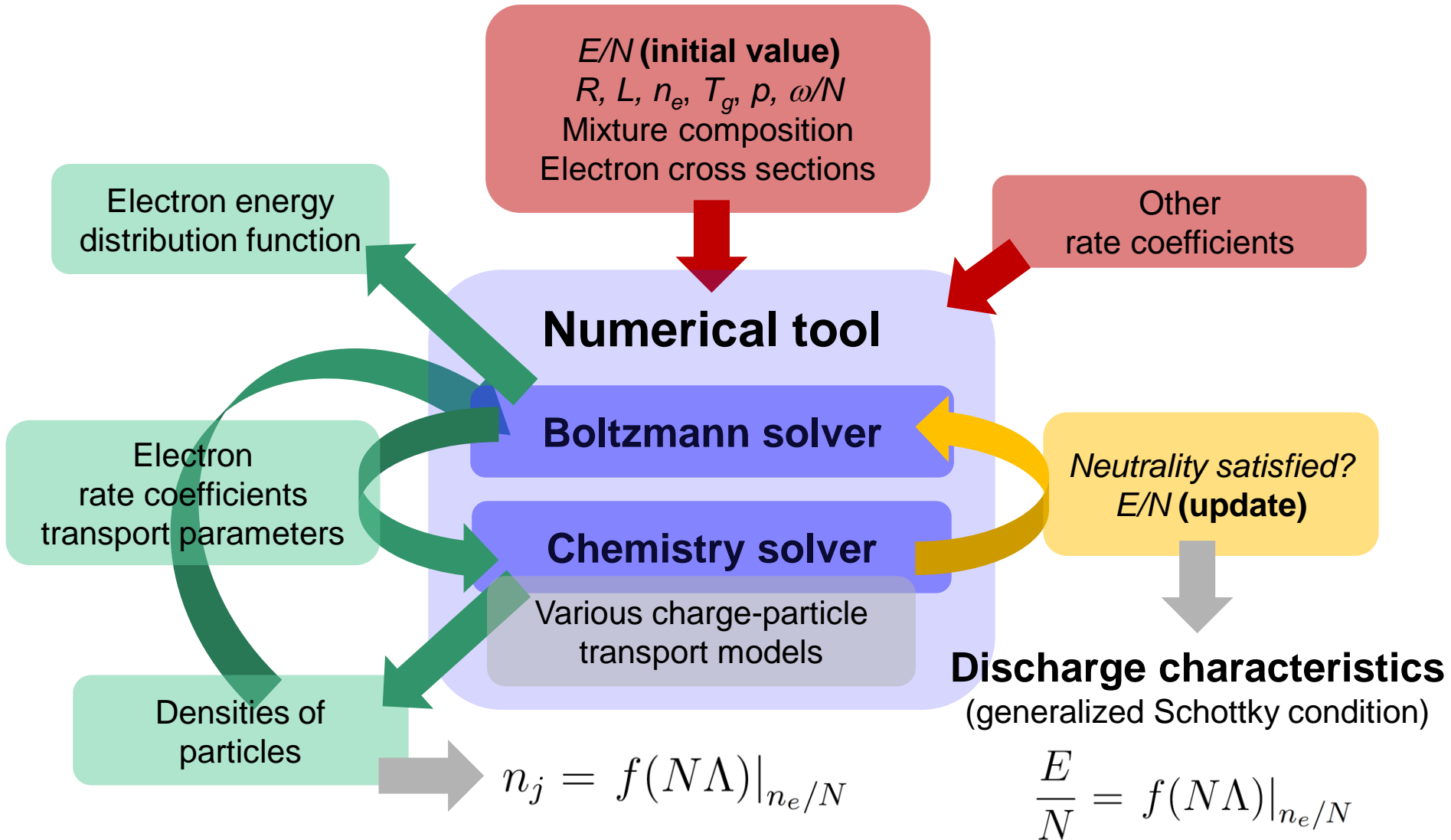
- solves the space independent form of the two-term electron Boltzmann equation, for DC/HF or time-dependent (non-oscillatory) electric fields.
- includes e-e collisions, CAR operator, and growth models for the electron density.

## LoKI-C

- solves the system of 0D rate balance equations for the heavy particles.
- includes modules to describe
  - the collisional, radiative and transport mechanisms controlling the creation / destruction of species
  - the thermal heating of the neutral gas



# LoKI workflow



## Cylindrical DC glow discharge in oxygen

Dias *et al.*, PSST (2023); Booth *et al.*, PSST (2019, 2020, 2022)

$R = 1 \text{ cm}$ ,  $L = 52.5 \text{ cm}$

$p = 25 - 1300 \text{ Pa}$  ;  $T_g \sim 320 - 580 \text{ K}$

$I_{dc} = 10 - 40 \text{ mA}$  ( $n_e \sim 10^{15} - 8 \times 10^{15} \text{ m}^{-3}$ )

Reaction mechanism considering

$O_2(X, v=0-41)$ ,  $O(^3P)$ ,  $O_3(X)$ ,  $O_2(a^1\Delta)$ ,  $O_2(b^1\Sigma_g)$ ,  $O_2(A^3\Sigma_u + C^3\Delta_u + c^1\Sigma_u)$ ,  $O(^1D)$ ,  $O_3^*$ ,  $O_2^+$ ,  $O^+$ ,  $O^-$

Wall recombination probability  $\gamma[O(^3P)] = f(p, T_g, I_{dc})$  from experiment

## Cylindrical MW discharge in helium

Alves *et al.*, JPD (1992); Santos *et al.*, PSST (2014); Alves and Tejero-del-Caz, PSST (accepted, 2023)

$R = 0.3 \text{ cm} \ll L$

$p = 50 - 5000 \text{ Pa}$  and  $10^5 \text{ Pa}$  ;  $T_g \sim 1000 - 1800 \text{ K}$

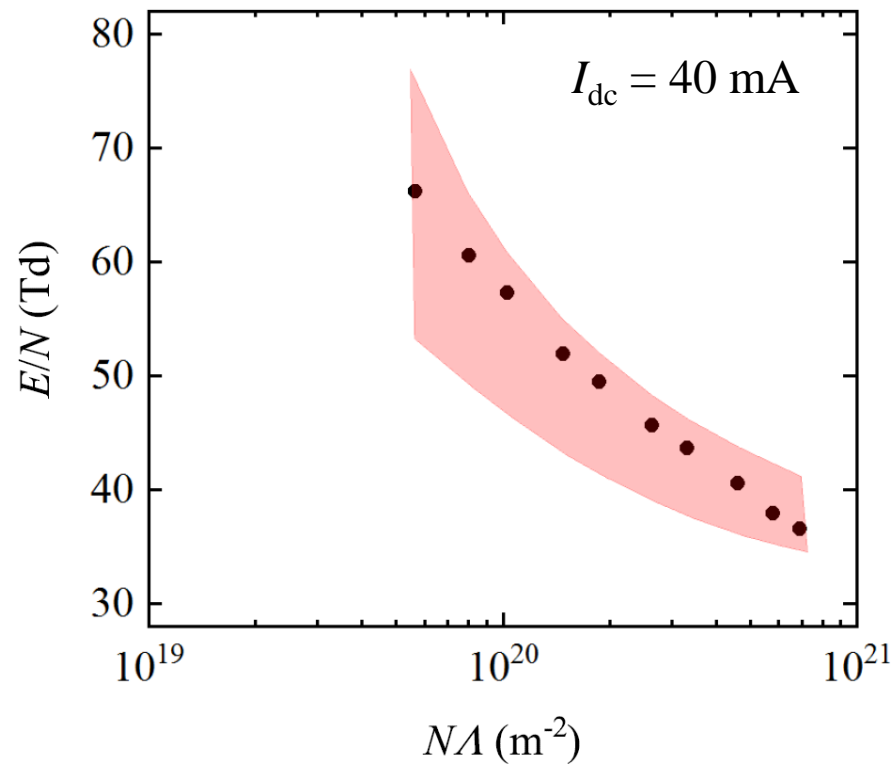
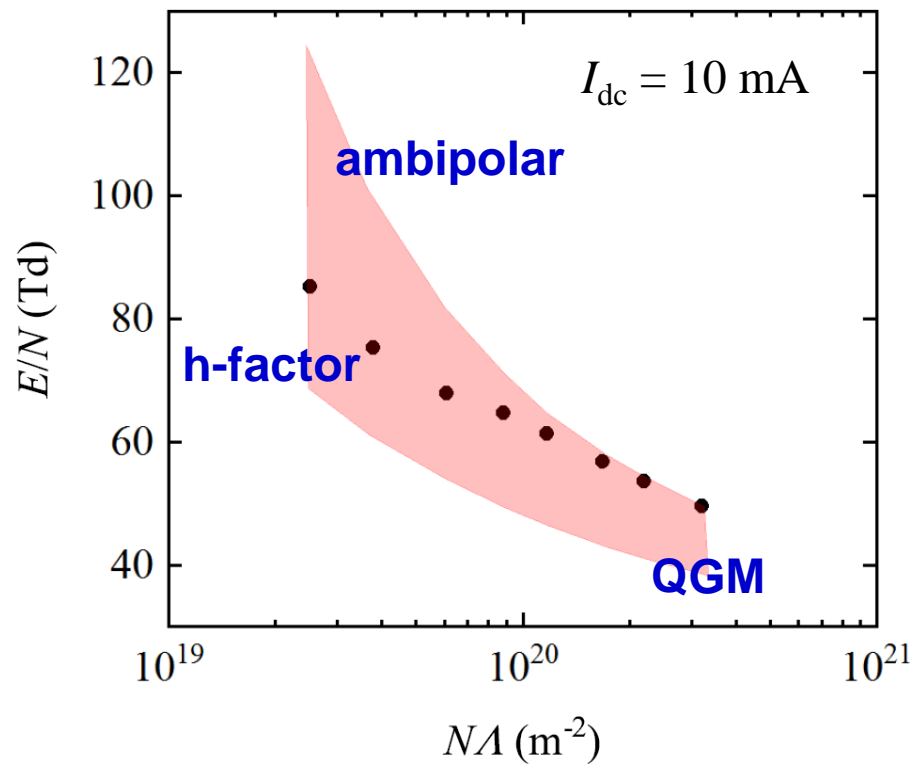
$\omega/(2\pi) = 2.45 \text{ GHz}$  ;  $n_e \sim 5 \times 10^{17} - 2 \times 10^{19} \text{ m}^{-3}$

Reaction mechanism considering

$He(n^{2s+1}l)$  ,  $n \leq 7$  ;  $He^+$ ,  $He_2^+$ ,  $He_2^*$

Extensive revision of AI rate coefficients  $He(1^1S) + He(n^{2s+1}l) \rightarrow e + He_2^+$

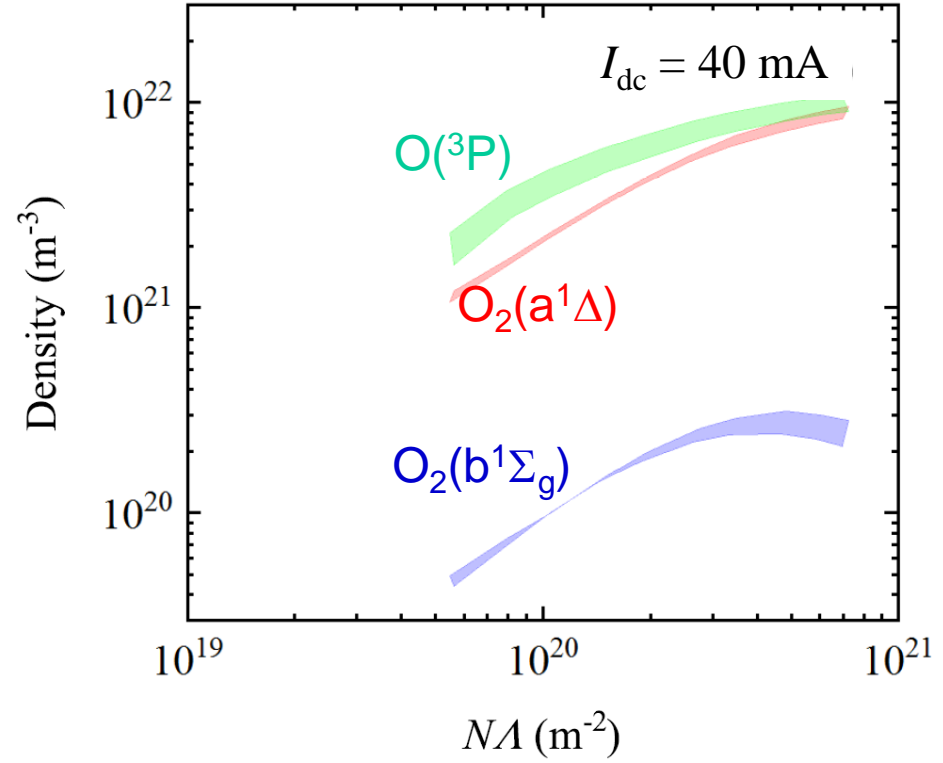
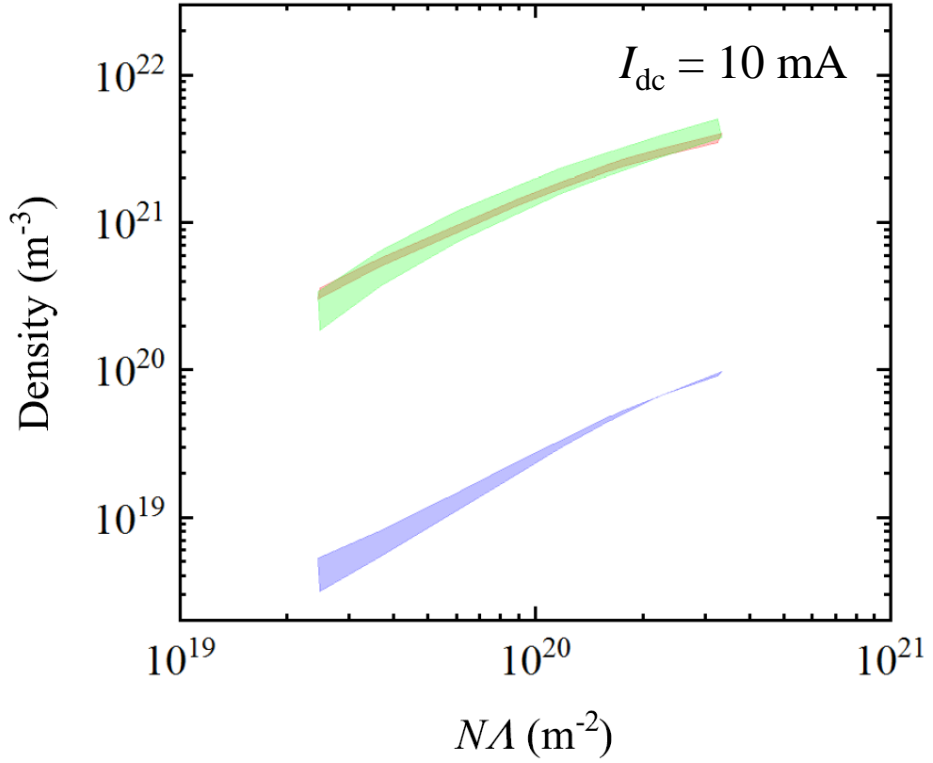
# Results in oxygen discharge characteristics



uncertainties of 20-60%  
larger dispersion at low pressure and low discharge current

# Results in oxygen

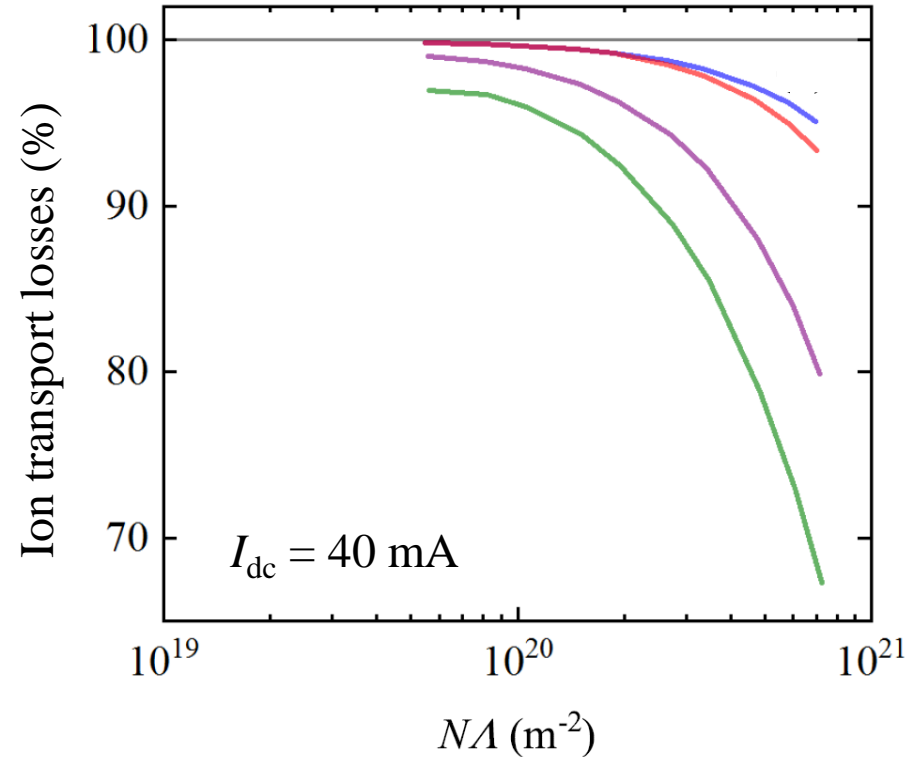
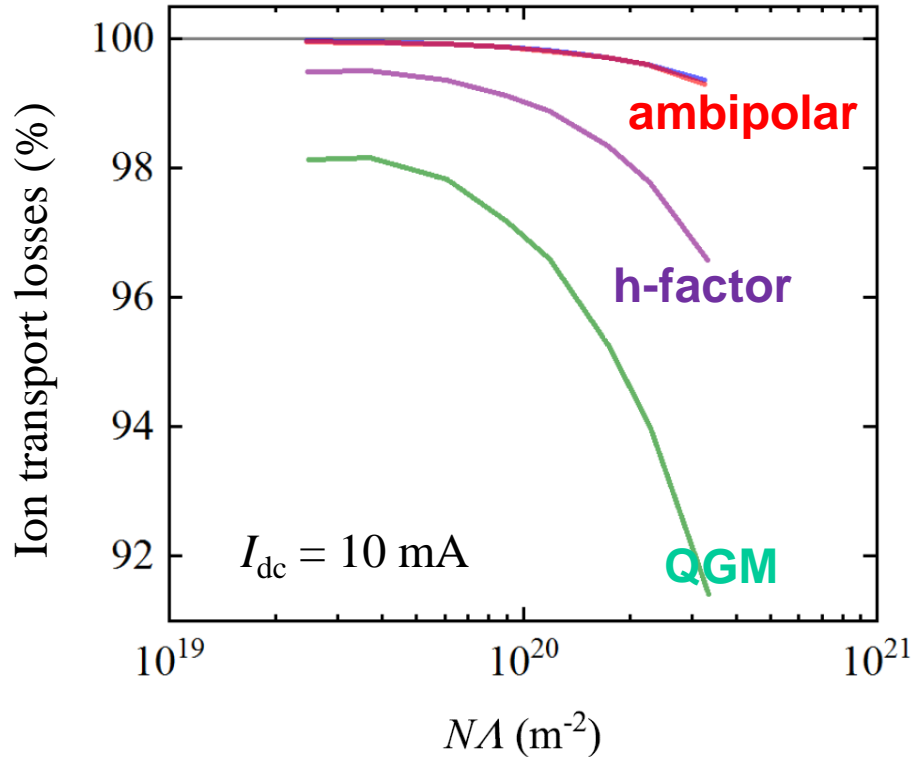
## densities of main plasma species



uncertainties that can reach 60%  
larger dispersion at low pressure and low discharge current

# Results in oxygen

## relevance of transport losses in the ion kinetics

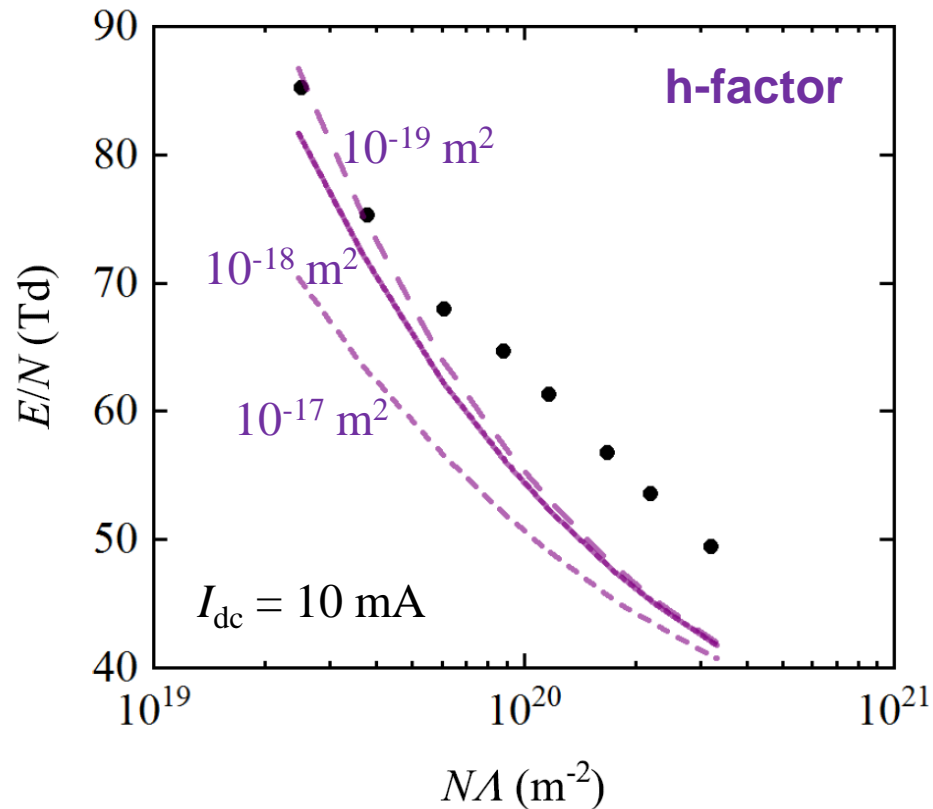


Transport accounts for > 70 – 90% of the total ion losses

- for  $O_2^+$ , transport is responsible for > 90% of destruction
- for  $O^+$ , charge-transfer  $O^+ + O_2(X, v=0) \rightarrow O_2^+ + O(^3P)$  yields 40-95% destruction

# Results in oxygen

influence of elementary data:  $\sigma_{ij}$  in h-factor model



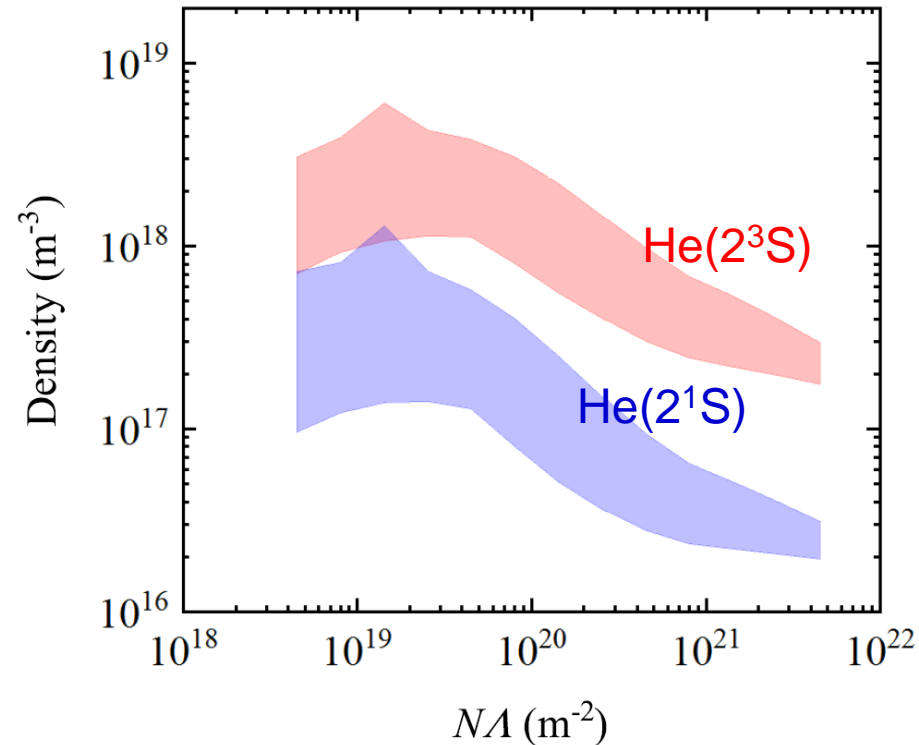
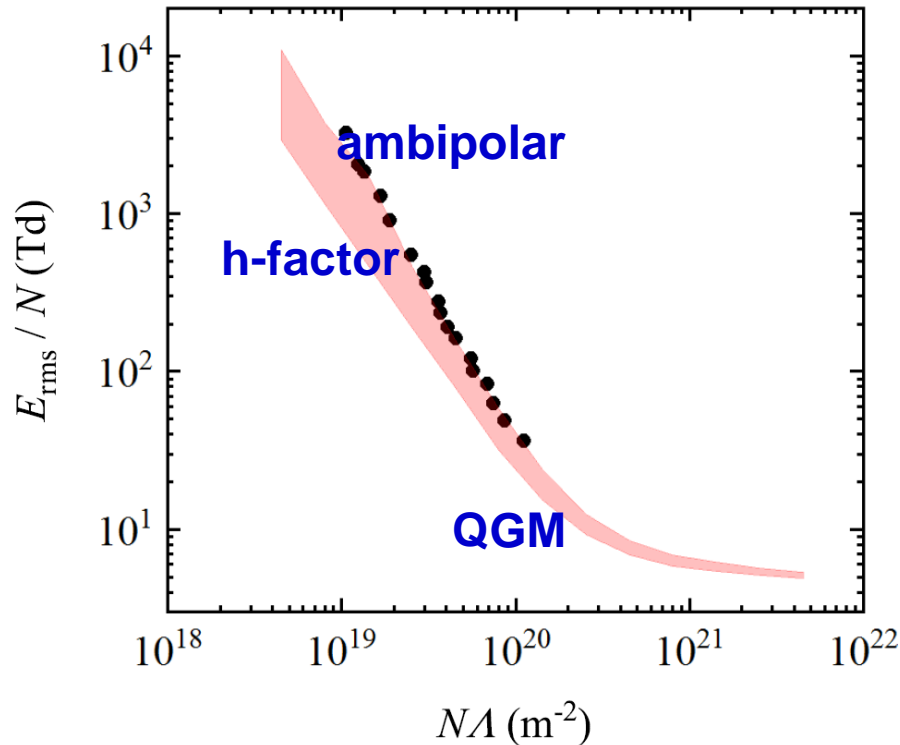
dispersion between 2 – 20%

the changes in the collisional cross section

- modify the LP terms only (dependent on  $\lambda_i$ )
- keep HP terms unchanged (dependent on  $D_i$ )

# Results in helium

## discharge characteristics and densities of main plasma species



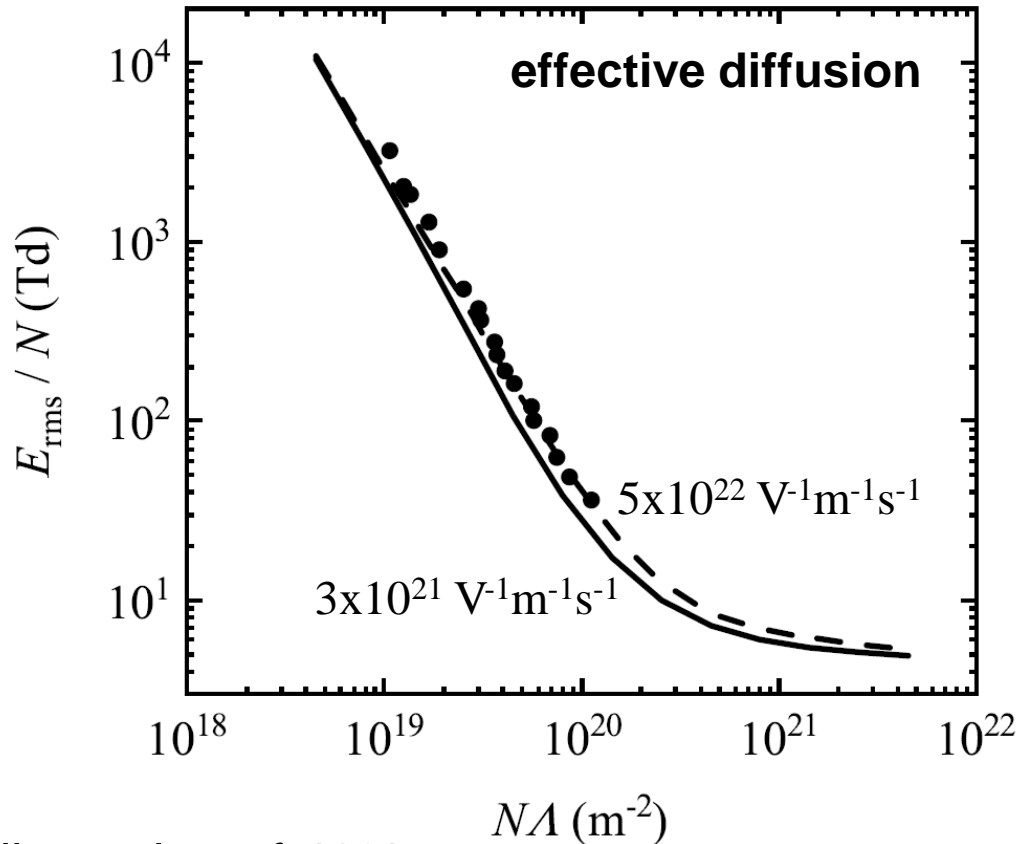
uncertainties of

- 8-115% for the discharge characteristics
- 50-150% for the densities of metastables

larger dispersion at low pressure

# Results in helium

influence of elementary data:  $\mu(\text{He}_2^+)N$  in effective diffusion model



The literature reports a dispersion  $> 25\%$  in  $\mu(\text{He}_2^+)N$

dispersion of 40%

the discharge characteristics become unchanged

- at LP, where transport becomes the dominant loss mechanism
- at HP, where transport losses are strongly reduced



# Final remarks

- We have analyzed the influence of different charged-particle transport models on the global modeling of LTPs
- The models were applied to the description of
  - a DC discharge in oxygen, at  $p \sim 25 - 1300$  Pa
  - a microwave discharge in helium, from  $p \sim 5 - 10^5$  Pa
- The simulations reveal dispersions of
  - 5 – 120% for the discharge characteristics
  - 50 – 150% for the densities of the main plasma species
- Transport accounts for more than 50% of total charge losses

**Charge-particle transport in global models is relevant**

# Final remarks

- **Transport phenomena should be considered and assessed in plasma chemistry studies** like any other kinetic mechanism, for example as part of the quantitative sensitivity analysis of a kinetic scheme
- To this date, **there is no general formulation** describing the transport of charged-particles in plasmas that can be used in global models for all types of plasmas and working conditions (this probably out of reach)
- The **choice of “suited” transport models** should consider the specific conditions considered (pressure, electronegativity, single-/multi-ions)
- Efforts should be made
  - to **translate accurate fluid/PIC-MCC simulation results into fitting / abacus formula** for transport loss-frequencies
  - to **express kinetic coefficients** (mobilities, diffusion coefficients, loss frequencies) as a **function of measured / calculated elementary data** (ion-scattering cross sections)

# Acknowledgements

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A Tejero-Del-Caz



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# FCT

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Recuperación,  
Transformación  
y Resiliencia